

PROBLEM SET 2

DUE DATE: FEB 28

- **Sections 2.4-3.5**

- Questions are either directly from the text or a small variation of a problem in the text.
- Collaboration is okay, but final submission must be written individually. Mention all collaborators on your submission.
- The terms in the bracket indicate the problem number from the text.

Section 2.4

1) (Prob 9, Pg 53) Solve the diffusion equation $u_t = ku_{xx}$ with the initial condition $u(x, 0) = x^2$ by the following method.

i) Show that u_{xxx} also satisfies the diffusion equation with zero initial condition.

Solution:

$$\begin{aligned}u_t &= ku_{xx} \\ \partial_{xxx}(u_t) &= k\partial_{xxx}u_{xx} \\ \partial_t(u_{xxx}) &= k\partial_{xx}(u_{xxx}) \quad (\text{Since partials commute})\end{aligned}$$

ii) Conclude that $u_{xxx}(x, t) \equiv 0$ for all (x, t) assuming $|u(x, t)| \leq Me^{ax^2}$ for all x, t . (Refer to practice problem set 2)

Solution: $u_{xxx}(x, 0) \equiv 0$ since $u(x, 0) = x^2$. Since under the growth conditions, we have uniqueness for the heat equation on the real line, we conclude that $u_{xxx}(x, t) = 0$

iii) Using ii, show that

$$u(x, t) = A(t)x^2 + B(t)x + C(t)$$

Solution: Follows from integrating the above expression 3 times.

iv) Plug in the above expression for $u(x, t)$ in the differential equation, to obtain a system of differential equations for $A(t), B(t), C(t)$. (Hint: the functions $1, x, x^2$ are linearly independent)

$$A'(t)x^2 + B'(t)x + C'(t) = \partial_t u = ku_{xx} = 2kA(t)$$

$$A'(t) = 0$$

$$B'(t) = 0$$

$$C'(t) = 2kA(t)$$

Solving the above system, we get

$$A(t) = a$$

$$B(t) = b$$

$$C(t) = 2kat + c$$

v) Using the initial conditions, obtain initial values for $A(0), B(0)$ and $C(0)$ and solve the above system of differential equations to compute $u(x, t)$

$$A(0)x^2 + B(0)x + C(0) = x^2$$

$$A(0) = 1, B(0) = 0, C(0) = 0.$$

$$u(x, t) = x^2 + 2kt$$

2) (Prob 15, Pg 53) Using the energy method, prove uniqueness of the diffusion problem with Neumann boundary conditions:

$$\begin{aligned}u_t - ku_{xx} &= f(x, t) \quad 0 < x < 1, t > 0 \\u(x, 0) &= \phi(x) \\u_x(0, t) &= g(t) \\u_x(1, t) &= h(t)\end{aligned}$$

Solution: To prove uniqueness for the Neumann problem, we consider the problem above with $f(x, t) \equiv 0$, $\phi(x) \equiv 0$, $g(t) \equiv 0$ and $h(t) \equiv 0$ and show that the solution

$$u(x, t) \equiv 0.$$

$$\begin{aligned}e(t) &= \int_0^1 u^2(x, t) dx \\e'(t) &= \int_0^1 2uu_t dx \\&= 2k \int_0^1 uu_{xx} dx \\&= 2k[uu_x]_0^1 - 2k \int_0^1 u_x^2 dx \quad (\text{Integration by parts}) \\&= -2k \int_0^1 u_x^2 dx \leq 0 \quad (u_x(0, t) = u_x(1, t) = 0)\end{aligned}$$

Thus, the energy is a decreasing function of time. Moreover $e(0) = 0$ since $u(x, 0) \equiv 0$. $e(t)$ is a non negative, decreasing function of time, which is 0 at $t = 0$, so we conclude that $e(t) \equiv 0$ and thus $u(x, t) \equiv 0$.

3) (Prob 16, Pg 54) Solve the diffusion equation with constant dissipation:

$$\begin{aligned}u_t - ku_{xx} + bu &= 0 \quad -\infty < x < \infty \\u(x, 0) &= \phi(x),\end{aligned}$$

where $b > 0$ is a constant by setting $u(x, t) = e^{-bt}v(x, t)$.

Solution:

$$\begin{aligned}v(x, t) &= e^{bt}u(x, t) \\ \partial_t v(x, t) &= e^{bt}\partial_t u(x, t) + be^{bt}u(x, t) \\ \partial_{xx} v(x, t) &= e^{bt}\partial_{xx} u(x, t) \\ \partial_t v - k\partial_{xx} v &= e^{bt}(\partial_t u + bu - k\partial_{xx} u) = 0\end{aligned}$$

Thus, $v(x, t)$ satisfies the heat equation with initial data $v(x, 0) = e^{b \cdot 0}\phi(x) = \phi(x)$. The solution for $v(x, t)$ is given by

$$\begin{aligned}v(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy \\u(x, t) &= e^{-bt} \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy\end{aligned}$$

Section 3.1

4) (Prob 2, Pg 60) Solve $u_t = ku_{xx}$; $u(x, 0) = 0$; $u(0, t) = 1$ on the half line $0 < x < \infty$. (Hint: Set $v(x, t) = u(x, t) - 1$.)

Solution: Let $v(x, t) = u(x, t) - 1$. Then v satisfies the differential equation

$$\begin{aligned}\partial_t v &= k\partial_{xx} v \\v(0, t) &= 0 \quad \forall t > 0 \\v(x, 0) &= -1 \quad 0 < x < \infty\end{aligned}$$

The solution is given by

$$\begin{aligned} v(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_0^\infty \left(\exp\left(-\frac{(x-y)^2}{4kt}\right) - \exp\left(-\frac{(x+y)^2}{4kt}\right) \right) \phi(y) dy \\ &= -\frac{1}{\sqrt{4\pi kt}} \int_0^\infty \left(\exp\left(-\frac{(x-y)^2}{4kt}\right) - \exp\left(-\frac{(x+y)^2}{4kt}\right) \right) dy \\ &= -\operatorname{erf}\left(\frac{x}{\sqrt{4kt}}\right) \\ u(x, t) &= 1 - \operatorname{erf}\left(\frac{x}{\sqrt{4kt}}\right) \end{aligned}$$

Section 3.2

5) (Prob 3, Pg 66) A wave $f(x + ct)$ travels along a semi-infinite string $0 < x < \infty$ for $t > 0$. Find the solution $u(x, t)$ of the string for $t > 0$ if the end $x = 0$ is fixed, i.e. $u(0, t) = 0$. Plot the solution in the three different regimes. Repeat the same exercise if $u_x(0, t) = 0$. Comment on the results.

Solution:

Dirichlet case

$$u(x, t) = \begin{cases} f(x + ct) & x > ct \\ f(x + ct) - f(ct - x) & x < ct \end{cases}$$

Neumann case

$$u(x, t) = \begin{cases} f(x + ct) & x > ct \\ f(x + ct) + f(ct - x) & x < ct \end{cases}$$

6) (Prob 5, Pg 66) Solve $u_{tt} = 4u_{xx}$ for $0 < x < \infty$, $u(0, t) = 0$, $u(x, 0) = 1$ and $u_t(x, 0) = 0$ using the reflection method. Find the location of the singularity of the solution in the (x, t) space.

Solution:

$$u(x, t) = \begin{cases} 1 & x \geq 2t \\ 0 & x < 2t \end{cases}$$

The singularity is on the line $x = 2t$.

Section 3.3

7) (Prob 2, Pg 71) Solve the completely inhomogeneous diffusion problem on the half line

$$\begin{aligned} v_t - kv_{xx} &= f(x, t) \quad 0 < x < \infty, \quad 0 < t < \infty \\ v(0, t) &= h(t), \quad v(x, 0) = \phi(x), \end{aligned}$$

by setting $V(x, t) = v(x, t) - h(t)$.

Solution: Let $V(x, t) = v(x, t) - h(t)$. Then $V(x, t)$ satisfies

$$\begin{aligned} V_t - kV_{xx} &= v_t - h'(t) - kv_{xx} = f(x, t) - h'(t) \\ V(0, t) &= 0 \\ V(x, 0) &= \phi(x) - h(0) \end{aligned}$$

By Duhamel's principle, the solution to the above inhomogeneous problem is given by

$$V(x, t) = \int_0^t \int_{-\infty}^{\infty} [S(x-y, t-s) - S(x+y, t-s)] \cdot [f(y, s) - h'(s)] dy ds + \int_{-\infty}^{\infty} [S(x-y, t) - S(x+y, t)] \cdot [\phi(y) - h(0)] dy$$

$$v(x, t) = h(t) + V(x, t)$$

Section 3.4

8) (Prob 12,13, Pg 80) Derive the solution of the fully inhomogeneous wave equation on the half-line

$$v_{tt} - c^2 v_{xx} = f(x, t), \quad 0 < x < \infty$$

$$v(x, 0) = \phi(x) \quad v_t(x, 0) = \psi(x)$$

$$v(0, t) = h(t)$$

by setting $V(x, t) = v(x, t) - h(t)$. Find the solution for $h(t) = t^2$, $\phi(x) = x$ and $\psi(x) = 0$.

Solution: Let $V(x, t) = v(x, t) - h(t)$. Then $V(x, t)$ satisfies the differential equation

$$V_{tt} - c^2 V_{xx} = f(x, t) - h''(t)$$

$$V(x, 0) = \phi(x) - h(0)$$

$$V_t(x, 0) = \psi(x) - h'(0)$$

$$V(0, t) = 0.$$

Let

$$\bar{\phi}_{odd}(x) = \begin{cases} \phi(x) - h(0) & x \geq 0 \\ -(\phi(-x) - h(0)) & x < 0 \end{cases}$$

$$\bar{\psi}_{odd}(x) = \begin{cases} \psi(x) - h'(0) & x \geq 0 \\ -(\psi(-x) - h'(0)) & x < 0 \end{cases}$$

$$\bar{f}_{odd}(x, t) = \begin{cases} f(x, t) - h''(t) & x \geq 0, t > 0 \\ -(f(-x, t) - h''(t)) & x < 0, t > 0 \end{cases}$$

By Duhamel's principle, the solution is given by

$$V(x, t) = \frac{1}{2} [\bar{\phi}_{odd}(x+ct) + \bar{\phi}_{odd}(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \bar{\psi}_{odd}(y) dy + \frac{1}{2c} \int_0^t \int_{x-c(t-s)}^{x+c(t-s)} \bar{f}_{odd}(y, s) dy ds$$

$$v(x, t) = V(x, t) + h(t)$$

$$v(x, t) = \begin{cases} x & x \geq ct \\ x + (t - \frac{x}{c})^2 & 0 \leq x \leq ct \end{cases}$$

9) (Stability to small perturbations for the heat equation) Consider the inhomogeneous heat equation on the real line

$$u_t - ku_{xx} = f(x, t) \quad -\infty < x < \infty, \quad t > 0$$

$$u(x, 0) = \phi(x).$$

Show that the solutions are stable under small perturbations, i.e. Show that

$$\|u\|_{\mathbb{L}^\infty(\mathbb{R} \times [0, T])} = \sup_{x \in \mathbb{R}, t \in [0, T]} |u(x, t)| \leq T \|f\|_{\mathbb{L}^\infty(\mathbb{R} \times [0, T])} + \|\phi\|_{\mathbb{L}^\infty(\mathbb{R})}.$$

Use the formula for the solution and the fact that

$$\int_{-\infty}^{\infty} S(x, t) dx = 1 \quad \forall t > 0.$$

Comment on why the above result implies stability of solutions to the data for the heat equation.

Solution:

$$\begin{aligned}
u(x, t) &= \int_0^t \int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) dy ds + \int_{-\infty}^{\infty} S(x-y, t) \phi(y) dy \\
|u(x, t)| &= \left| \int_0^t \int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) dy ds + \int_{-\infty}^{\infty} S(x-y, t) \phi(y) dy \right| \\
&\leq \left| \int_0^t \int_{-\infty}^{\infty} S(x-y, t-s) f(y, s) dy ds \right| + \left| \int_{-\infty}^{\infty} S(x-y, t) \phi(y) dy \right| \quad (\text{Triangle ineq}) \\
&\leq \int_0^t \int_{-\infty}^{\infty} |S(x-y, t-s) f(y, s)| dy ds + \int_{-\infty}^{\infty} |S(x-y, t) \phi(y)| dy \quad \left(\left| \int f \right| \leq \int |f| \right) \\
&\leq \|f\|_{\mathbb{L}^\infty(\mathbb{R} \times [0, T])} \int_0^t \int_{-\infty}^{\infty} S(x-y, t-s) dy ds + \|\phi\|_{\mathbb{L}^\infty(\mathbb{R})} \int_{-\infty}^{\infty} S(x-y, t) \phi(y) dy \quad (S > 0) \\
&\leq \|f\|_{\mathbb{L}^\infty(\mathbb{R} \times [0, T])} \int_0^t ds + \|\phi\|_{\mathbb{L}^\infty(\mathbb{R})} \left(\int_{\mathbb{R}} S dx = 1 \right) \\
&\leq T \|f\|_{\mathbb{L}^\infty(\mathbb{R} \times [0, T])} + \|\phi\|_{\mathbb{L}^\infty(\mathbb{R})}
\end{aligned}$$

The above result implies stability because small perturbations in f and ϕ , result in small perturbations in the solution u .